Outflows from young stellar objects and their impact on star formation

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Jets and outflows are observed around young stellar objects over the whole stellar spectrum, from brown dwarfs to high-mass stars. Those outflows are most likely driven by the coupling of magnetic fields that thread the underlying accretion disc. If this is a universal mechanism, such a disc-wind configuration should be self-consistently build up during the collapse of individual cloud cores. Additionally, jets and outflows feed back energy and momentum to the ambient gas in star-forming regions. Yet, it is still controversial whether feedback from outflows are able to regulate star formation in molecular clouds.

In this talk, I will summarise recent results from numerical simulations on outflow launching during the birth of stars and their impact on star-forming regions based on sub-grid modells of outflows.

Subject :	:	oral
Topics	:	Astrophysics
Topics	:	Numerical simulations



Outflows from YSOs and their Impact on Star Formation

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Collaborators: Christoph Federrath (Monash), Daniel Seifried (Cologne), Thomas Peters (MPA)

Outflows & Jets

- Outflows & Jets are ultimately linked to the formation of stars
 - ⇒ what's their impact on this process?
 - ⇒ how to model it self-consistently?



Collapse of Magnetised Cloud Cores





magnetically driven Jets / Outflow from YSOs

Onset of large scale outflow:

at few 100 AU

magnetic tower configuration (e.g. Lynden-Bell 2003)



collapse phase pinched in magnetic field

.... I 430 years later: onset of a large scale outflow

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Banerjee & Pudritz 2006

Magnetic tower flow



- build up of toroidal field \rightarrow magnetic pressure
- outward propagation of shock fronts
- magnetic bubble

Magnetic tower flow



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- outward propagation of shock fronts
- magnetic bubble

Large scale outflow



- Magnetic field is
 compressed with the gas
- Rotating disk generates
 toroidal magnetic field
 ⇒outflow
- Shock fronts are pushed outwards (magnetic tower)
- •Outflow velocities v ~ 0.4 km/sec, M ~ 2-3
- •Accretion: funneled along the
 - rotation axis, through disk

$\Delta x \approx 5 \times 10^9 \text{ cm} (0.07 \text{ R}_{sol}) \text{ at } l_{ref} = 27$ $\implies \text{Onset of inner disk jet}$



small scale disk jet



- Magnetic field strongly pinched and warped
- •Angle with disk plane < 60°
- → magneto-centrifugal jet launch (Blandford & Payne 1982)
- "Onion" shaped velocity structure
- Outflow velocities
 - v ~ 4 5 km/sec, Mach ~ 4

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 \implies can not follow long-term, large-scale evolution

Sink Particles for the FLASH code

 AMR/SPH simulations can't cover the full spatial range for star formation => introduce "black boxes" = Sink Particles

Sink Particles for the FLASH code

- AMR/SPH simulations can't cover the full spatial range for star formation => introduce "black boxes" = Sink Particles
 - \implies modeling of dense regions in **collapse** simulations,

e.g. star formation (M.Bate et al. 1995)

 'controlled' violation of the Truelove criterion (*Truelove et al. 1997*): preventing artificial fragmentation by resolving the

Jeans length

- allows long term runs of star forming regions: binaries, stellar clusters, outflows also: feedback, drag forces, ...
- **BUT:** 'arithmetic' part of the system, i.e.
 ⇒ physical interpretation?



Based on the particle module in FLASH 2.x (Paul Ricker):

- handles boundaries
- moves particles across CPUs/blocks
- mapping of grid variables onto particles and vice versa
- advances particles

Extensions / modifications:

- creation of particles on the 'fly'
- gravity: use $1/r^2$ acceleration for particle contribution
- time dependent particle masses: accretion / loss
- momentum transfer onto the particles
- back-reaction onto the grid (feedback)
- MPI communication of global particle list

Gravity

originally:

- I. mapping of particle density onto grid (CIC, NGP, TSC)
- 2. solve Poisson's equation with gas-density + particle-density
- 3. map acceleration to particle
- 4. advance particle (Euler, Leapfrog)

Sink Particle Module:

- use direct acceleration from particles
 - ⇒ more accurate (e.g. binary system)
 - \Rightarrow faster for small particle numbers

• Gravity



with $1/r^2 \rightarrow f(r, r_{soft})$: gravitational softening

Gravitational softening



- Sub-Cycling
 - close "binary" interaction can limit time step:

$$\Delta t_{\rm gs} = C_{\rm gs} \, \min_{n,m} \left(\frac{\min(|\mathbf{r}_{nm}|, \Delta x)}{|\mathbf{g}_{\rm sinks, n}|} \right)^{1/2}$$

 \Rightarrow sub-cycle on particle-particle interaction till:

$$N_{\rm cycles} \,\Delta t_{\rm gs} = \Delta t_{\rm hydro}$$

Sub-Cycling



\Rightarrow after 10 orbits

Sub-Cycling



 \Rightarrow after 1000 orbits:

two particles around the common center

Particle creation

Conditions by gravitational **collapse**:

0. Density criterion (within accretion radius r_{accr}):

$$\rho_{\rm gas} > \rho_{\rm crit} \quad (\rho_{\rm crit} \text{ parameter})$$

 \Rightarrow choose ρ_{crit} so that Truelove criterion is not violated:

$$\lambda_{\rm J} > N_{\rm J} \Delta x_{\rm min}$$

+ Jeans refinement condition ($\lambda J = (\pi c 2/G\rho)^{1/2}$)

Particle creation

Conditions by gravitational collapse:

0. Density criterion: $\rho_{gas} > \rho_{crit}$ (ρ_{crit} parameter)

- 1. is on the highest level of refinement,
- 2. is converging, $\nabla \cdot \mathbf{v} < 0$
- 3. has a central gravitational potential minimum,
- 4. is Jeans-unstable, $|E_{grav}| > 2E_{th}$
- 5. is bound, and $E_{\text{grav}} + E_{\text{th}} + E_{\text{kin}} + E_{\text{mag}} < 0$
- 6. is not within r_{acc} of an existing sink particle.

Federrath, RB, Clark & Klessen et al. 2010

Sink Particles in FLASH

Particle creation



Sink Particles in FLASH

Mass accretion & linear momentum

Mass accretion from excess gas density within $r_i < r_{accr}$:

$$\mathbf{M}_{i} = \mathbf{M}_{i} + \Sigma_{j} \Delta \operatorname{Vol}_{j} \left(\rho_{j} - \rho_{\operatorname{crit}} \right)$$

additional check for convergent flow, i.e. $v_{\rm rad} < 0$ Mass conservation ensured

+ linear momentum conservation:

$$\mathbf{P}_i = \mathbf{P}_i + \Sigma_j \Delta \mathbf{m}_j \mathbf{v}_j$$

angular momentum

no unique solution for angular momentum conservation:

$$\mathbf{R} \times \mathbf{v}_{\rm cm} = \frac{1}{M} \mathbf{L}$$

 \Rightarrow internal spin:

$$\mathbf{L}_{\mathrm{spin}} = \mathbf{L}_{\mathrm{gas}}' - \mathbf{L}_{\mathrm{gas}}$$

use for sub-grid-scale modelling, e.g. outflows & jets

Comparison to SPH Simulations

Federrath et al. 2010

• collapse of **turbulent** cloud cores

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Comparison to SPH Simulations



Federrath et al. 2010

• collapse of **turbulent** cloud cores

Comparison to SPH Simulations



- good agreement
- differences due to hydro
 - \Rightarrow SPH slightly more dissipative
 - \Rightarrow cluster more centrally condensed

Applications





disc formation and jet launching by Daniel Seifried



feedback from ionizing radiation by *Thomas Peters*

Outflows from Massive Stars: Young HII Regions

3D Simulations of collapsing cloud cores with ionization feedback from young massive stars (*Thomas Peters*, ITA)

- massive core with $M_{\rm core} = 1000 \ M_{\odot}$
- $R_{core} = 1.6 \text{ pc}$
- $\rho_{core} = 1.27 \times 10^{-20} \text{ g cm}^{-3}; \rho \sim r^{-1.5}$
- initial core rotation with $\beta = 0.05$
- magnetized case: $\mu = 14 \mu_{crit} (B = 10 \mu G)$



- accreting sink particles ⇒ luminosity and temperature using ZAMS (*Paxton* 2004)
 + protostellar accretion luminosity (*Hosokawa & Omukai* 2009)
- highest grid resolution $\sim 100 \text{ AU}$
- ray-tracing based on Rijkhorst et al. 2006

Massive Star Formation: Dynamics of HII Regions

Run B: formation of multiple stars



courtesy: Zilken, NIC, Jülich

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Massive Star Formation: Dynamics of HII Regions

Collapse of a massive, rotating cloud core $(M_{core} = 1000 M_{sol})$ + ionization feedback

Simulations by Thomas Peters



Disk edge on

Disk plane

Massive Star Formation: Dynamics of HII Regions



- Ionization feedback does not shut off star formation
- accretion onto the most massive star is cut off by **fragmentation induced starvation** (Peters et al. 2010)

Comparison with Observations: Outflows



Synthetic CO maps with the ALMA simulator CASA
 Orion distance: 414 pc

Comparison with Observations: Outflows

	OUTFLOW PARAMETERS DERIVED FROM ALMA SIMULATIONS										
		М	v	Р	Е	L	Ń	Т	R		
		(M_{\odot})	$(\mathrm{km}\ \mathrm{s}^{-1})$	$(M_{\odot} \text{ km s}^{-1})$	(10^{44} erg)	(L_{\odot})	$(10^{-3} \ { m M}_{\odot} \ { m yr}^{-1})$	(yr)	(AU)		
Run A	blue	$2.50{\pm}0.26$	$3.9 {\pm} 0.9$	$9.93 {\pm} 3.32$	3.94 ± 2.22	$8.15 {\pm} 6.23$	6.26 ± 1.91	400	4100		
	red	$1.80 {\pm} 0.18$	$3.8{\pm}0.9$	$6.82 {\pm} 2.30$	$2.58{\pm}1.48$	$5.33{\pm}4.12$	$4.51{\pm}1.37$	400	4100		
Run B (left)	blue	$1.12{\pm}0.13$	3.3 ± 0.4	$3.68 {\pm} 0.87$	1.21 ± 0.43	2.51 ± 1.39	$2.80 {\pm} 0.89$	400	3300		
27 20	red	$2.08{\pm}0.12$	$3.5{\pm}0.5$	$7.26{\pm}1.55$	$2.53{\pm}0.93$	$5.24 {\pm} 2.98$	$5.21{\pm}1.35$	400	2100		
Run B (right)	blue	$1.31 {\pm} 0.12$	$3.3 {\pm} 0.4$	4.29 ± 0.91	1.41 ± 0.47	$2.92{\pm}1.55$	$3.26{\pm}0.95$	400	5000		
	red	$0.75{\pm}0.08$	$3.5{\pm}0.5$	$2.62{\pm}0.69$	0.91 ± 0.38	$1.89{\pm}1.17$	$1.88{\pm}0.58$	400	4100		

TABLE 1 OUTFLOW PARAMETERS DERIVED FROM ALMA SIMULATIONS

Peters, Klaassen et al. 2012



 → derived outflow parameters are on the low end of observations
 → lonisation feedback is not the main driver of molecular outflows
 → common low mass companions drive large scale molecular outflows? (see also Peters et al. 2014, Collective Outflows ...)

Magnetic fields during Massive Star Formation?



outflows launched by magnetic fields?

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Magnetic fields during Massive Star Formation?



Ideal MHD equations + (self-)gravity

references: eg. Chandrasekhar 1956; Mestel 1969; Blandford&Payne 1982; Pudritz&Norman 1983 reviews: eg. Königl&Pudritz 1999 (PPIV); Heyvaerts 2000; Pudritz et al. 2007 (PPV)

$$\begin{aligned} \frac{\partial \rho}{\partial t} + \nabla \cdot (\mathbf{v}\rho) &= 0\\ \frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} &= -\frac{1}{\rho} \nabla p - \nabla \Phi + \frac{(\nabla \times \mathbf{B}) \times \mathbf{B}}{4\pi\rho}\\ \frac{\partial \mathbf{B}}{\partial t} &= \nabla \times (\mathbf{v} \times \mathbf{B})\\ \nabla \cdot \mathbf{B} &= 0\\ \Delta \Phi = 4\pi G\rho \end{aligned}$$

Lorentz force:

(assume axi-symmetry, i.e. $\partial_{\Phi} \mathbf{B} = 0$)

$$\mathbf{j} \times \mathbf{B} = -\frac{1}{2} \nabla \mathbf{B}^2 + (\mathbf{B}_{\mathrm{p}} \cdot \nabla) \left(\mathbf{B}_{\mathrm{p}} + B_{\phi} \mathbf{e}_{\phi} \right) \underbrace{-\frac{B_{\phi}^2}{R} \mathbf{e}_R}_{R}$$

hoop stress (jet collimation)

Lorentz force:

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$$\mathbf{j} \times \mathbf{B} = -\frac{1}{2} \nabla \mathbf{B}^2 + (\mathbf{B}_{\mathrm{p}} \cdot \nabla) \left(\mathbf{B}_{\mathrm{p}} + B_{\phi} \mathbf{e}_{\phi} \right) \underbrace{-\frac{B_{\phi}^2}{R} \mathbf{e}_R}_{\mathbf{p}}$$

hoop stress (jet collimation)

2

different force types:

- magnetic pressure: force along gradient
- tension: force along magnetic field lines
- hoop stress: force towards axis

Lorentz force:

(assume axi-symmetry, i.e. $\partial_{\Phi} \mathbf{B} = 0$)

$$\mathbf{j} \times \mathbf{B} = -\frac{1}{2} \nabla \mathbf{B}^2 + (\mathbf{B}_{\mathrm{p}} \cdot \nabla) \left(\mathbf{B}_{\mathrm{p}} + B_{\phi} \mathbf{e}_{\phi} \right) - \frac{B_{\phi}^2}{R} \mathbf{e}_R$$



0



"beads on a wire" Blandford-Payne type acceleration

magnetic pressure acceleration



courtesy Matsumoto & Shibata, 1999

Jet & Outflow Launching

Specific energy conserved along field lines

with separation of poloidal and toroidal velocity and field components:

$$\begin{aligned} \epsilon &= \frac{1}{2}v^2 + \Phi + h - \frac{r\omega B_{\phi}}{4\pi k} \\ &= \frac{1}{2}v_{\rm pol}^2 + \frac{1}{2}v_{\phi}^2 + \Phi + h - \frac{v_{\phi}}{v_{\rm pol}}\frac{1}{4\pi}\frac{B_{\phi}B_{\rm pol}}{\rho} + \frac{1}{4\pi}\frac{B_{\phi}^2}{\rho} \end{aligned}$$

 \implies new, generalised outflow criterion

to distinguish between tower and centrifugal launching

mew, generalised outflow criterion to distinguish between tower and centrifugal launching

⇒ magneto-centrifugal launching (a-la Blanford & Payne):

$$\frac{r}{z}\frac{1}{GM}\left(\frac{v_{\phi}^2}{r^2}(r^2+z^2)^{3/2}-GM\right)\left/\left(\frac{B_z}{B_r}\right)>1$$

 \implies any outward acceleration:

$$\partial_{\text{pol}} \left(\frac{1}{2} v_{\phi}^2 + \Phi - \frac{v_{\phi}}{v_{\text{pol}}} \frac{1}{4\pi} \frac{B_{\phi} B_{\text{pol}}}{\rho} + \frac{1}{4\pi} \frac{B_{\phi}^2}{\rho} \right) < 0$$

Seifried et al. 2012

Collapse of Massive Cloud Cores

Parameter study with 3D Simulations of massive collapsing cloud cores with Sink Particles

- $M_{core} = 100 M_{\odot}$
- $R_{core} = 0.125 \text{ pc}$
- density profile: $\rho \sim r^{-1.5}$
- $\rho_{core} = 2.3 \times 10^{-17} \text{ g cm}^{-3}$
- rotation with $\beta = 0.0004 0.2$
- mass-to-flux: $\mu = 2.6 \dots 26 \mu_{crit}$
- $B_z = 1.3 0.13$ mG aligned with rotation axis
- resolution: 4.7 AU



Seifried, RB, Klessen, Duffin, Pudritz 2011

A Generalised Outflow Criterion

Outflow / Lauchning mechanism



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A Generalised Outflow Criterion

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Parameter study of collapsing cores

Outflow / Launching mechanism

stronger magnetic field: $\mu = 5.2 \ \mu_{crit}$



- inefficient magneto-centrifugal launching
- bubble like "outflow"

Synthetic Observations



 \implies Helical structure similar to outflow around the A-type star HD 163296 (D = 122 pc)

Sub-Grid-Scale Model

SGS Model: Single Outflow

Low resolution No subgrid model High resolution No subgrid model Low resolution With SGS outflow model

Federroth et al. (2014)

SGS Model: Single Outflow

Low resolution No subgrid model High resolution No subgrid model

Low resolution With SGS outflow model

ederrath et al. (2014)

⇒ low resolution SGS outflow model recovers fast jet of high resolution self-consistent outflow simulation

• Jets are powerful:

$$L_{jet} = \frac{\dot{M}_{jet}v_{jet}^2}{2} \approx 2.9 \times 10^{32} \left(\frac{\dot{M}_{jet}}{10^{-8} M_{\odot} \text{ yr}^{-1}}\right)$$
$$\times \left(\frac{v_{jet}}{300 \text{ km s}^{-1}}\right)^2 \text{ ergs s}^{-1} \sim 8\% L_{\odot}$$
$$E_{jet} = L_{jet}\tau_{jet} \approx 10^{44} \text{ ergs} \qquad \text{with } \tau_{jet} = 10^4 \text{ yrs}$$
$$\Rightarrow \text{ cf. } E_{turb} \sim 10^{46} \text{ ergs}$$

 \implies Jets from a little stellar cluster **could** maintain the turbulence

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 \implies Jets from a little stellar cluster **could** maintain the turbulence

\implies But how **efficient** do they couple to the ISM?

- numerical experiments with single, high Mach number jets (momentum injection)
 detailed analysis with velocity PDFs
- log₁₀(velocity) log₁₀(velocity) 2 2 0 (-2 -2 10 8 10 2 6 2 8 4 4 6 t = 5.00t = 3.00RB, Klessen & Fendt 2007

mber of blocks =

- supersonic fluctuations
 decay quickly: E∝t⁻²
 (Mac Low et al. '98)
- supersonic fluctuations
 occupy only a small
 fraction of all fluctuations

Influence of Magnetic Fields

t = 2.00

magnetic fields **suppress** the propagation of large amplitude velocity fluctuations

stabilize jet (aligned field)

Influence of Magnetic Fields

Global simulation

• collapse of a turbulent

cloud core (Li&Nakamura 2006; Carroll et al. 2008, Dale & Bonnell 2008, Wang et al. 2010, Federrath et al. 2014)

Global simulation

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SGS Model: Outflows during Cluster Formation

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SGS Model: Outflows during Cluster Formation

Outflows & Jets do not stop star formation

Wang et al. (2010): Collapse of a massive, turbulent cloud core $(M_{core} = 1600 M_{sol}) + feedback$ from jets & outflows

Wang, Li, Abel & Nakamura 2010

 \implies Outflows & Jets do not stop star formation

Conclusion

- Jets & Outflows: self-consistent treatment in collapse simulation is still challenging (but see Hennebelle et al.)
- SGS models allow to scan a larger parameter space (at lower resolution)
- Influence of Outflow feedback?
 - \implies **not** conclusive:
 - \implies might not be too important on cloud scales

→ v= 11.4620 (km/s) Hennebelle et al. 2011