

# Radiative implications of obstacles embedded in microquasar and extragalactic jets

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The interaction of microquasar and extragalactic jets with relatively small obstacles, such as stellar wind inhomogeneities, clouds, or stars, can lead to important consequences from the radiative point of view. Particles can be accelerated in the interaction region up to very high energies either through Fermi mechanisms or magnetic reconnection, and produce Doppler boosted emission in different energy bands. In this talk, I will overview some important aspects of the non-thermal phenomena related to the entrainment of obstacles in galactic and extragalactic jets.

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Subject : : oral  
Topics : : Astrophysics

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Valentí Bosch-Ramon

Universitat de Barcelona/ICC

**Accretion and Outflows throughout the scales**

Lyon, October 3, 2014

- 1 Introduction
- 2 Radiation
- 3 Radiation and dynamics coupling
- 4 Conclusions

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## General context: jet environment

- 1 The ambient medium within massive binary systems and the nuclear regions of AGN is complex, rich in matter, magnetic and radiation fields.**
- 2 The presence of stars, or clumps of material originated by different types of processes (instabilities, winds, tidal forces, irradiation, stellar collisions...), render the environment of relativistic jets strongly inhomogeneous.

How much relevant can the medium inhomogeneities be for the jet dynamics and radiation?

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## Moving, dense obstacles can interact with jets

- 1 **Stellar wind inhomogeneities of different origin (MQ)**
- 2 Molecular clouds, clouds from the BLR, etc. (AGN)
- 3 Massive stars and red giants within galaxies (AGN)

Clouds can penetrate in the jet as long as

$$L_j/S_j c < \rho_0 v_0^2$$

(e.g. Blandford & Koenigl 1979; Coleman & Bicknell 1985; Steffen et al. 1997; Hubbard & Blackman 2006; Owocki et al. 2009)



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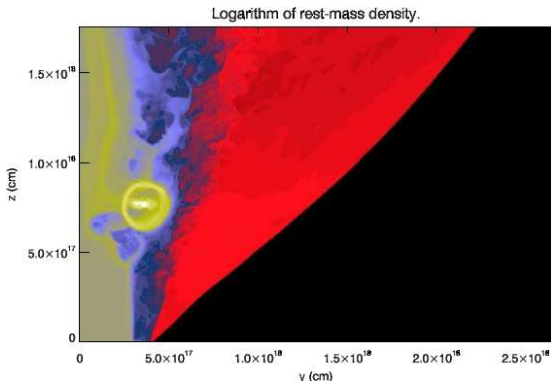
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A red giant entering in an AGN jet (density map).

(Perucho, B-R, Barkov, in prep.)

## These obstacles trigger instabilities affecting jet propagation

- 1 **Jets suffer kinetic energy dissipation: slow-down and heating**
- 2 Jets entrain matter: mass-load
- 3 Jets suffer strong instabilities: disruption

Stellar mass-load may be important decelerating/disrupting for FR I jets, but the stellar populations may not be well known.

The role of clouds may be also important.

Stellar wind clumps are likely dynamically relevant in high-mass microquasar jets.

(e.g. Komissarov et al. 1994, Bowman et al. 1996, Perucho & B-R 2012, B-R, Perucho & Barkov 2012, Laing & Bridle 2014, Perucho et al. 2014)

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## Jet-obstacle interaction can generate

- 1 **Strong shocks and velocity gradients**
- 2 Magnetic activity
- 3 Different emitting sites

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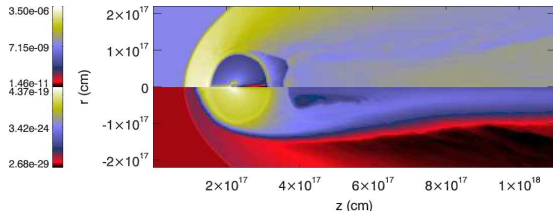
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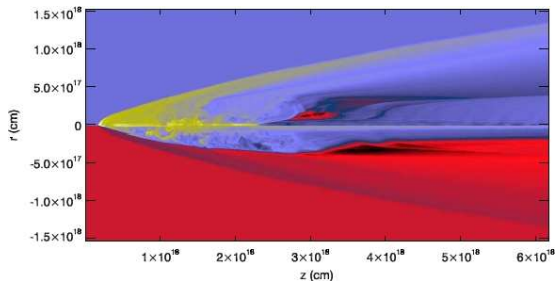
## RG bubble dragged by the jet (density map).

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- 1 **Magnetic pressure is much smaller than the momentum flux.**
- 2 The obstacle inertia stops the flow, which is supermagnetosonic.
- 3 A bow-shaped shock surrounds the obstacle.

Diffusive shock acceleration can accelerate particles if enough chaotic field is present. Shear acceleration may also accelerate particles far downstream.

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- 1 **If  $B_\phi$  is not negligible, magnetic energy can accumulate in front of the obstacle.**
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All this can trigger magnetic field reconnection, turbulence, and particle acceleration.

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- 1 **If  $B$  presents polarity reversals, it can reconnect in the accumulation region.**
- 2 Magnetic energy turns into heat (and kinetic energy far downstream) and potentially non-thermal particles.

Electric field acceleration can occur in the current sheet, and Fermi I acceleration can occur between the sheet and/or the in-coming flows.

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## Simulation: basic scenario

- 1 **Clumps in a stellar wind that can penetrate a microquasar jet on binary scales.**
- 2 Clouds can interact with an AGN jet on different scales.
- 3 RG and massive star winds can interact with an AGN jet on different scales.

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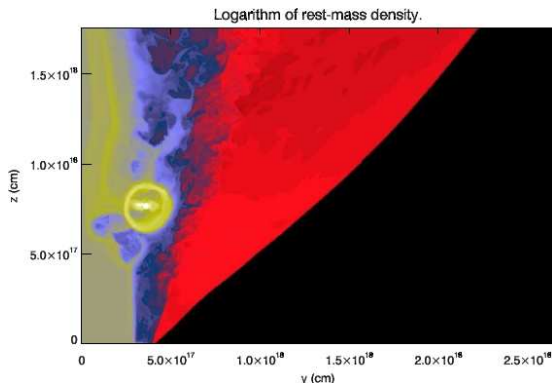
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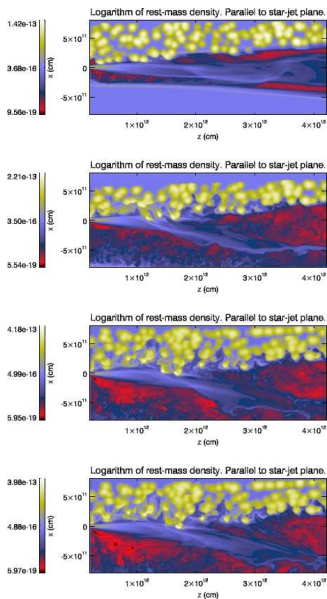
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Clumps entering into a MQ jet (density map):

(Perucho & B-R 2012)



## Simulation: assumptions

- 1 **Relativistic hydrodynamical flow.**
- 2 **Axisymmetric relativistic simulation.**  
(2nd spatial order, Marquina solver,  $300 \times 600$  cells;  $R_o = 30$  cells)
- 3 **Adiabatic interaction; relativistic ideal gas.**
- 4 **A jet meets an obstacle of very large inertia.**  
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- 2 The post-shock flow velocity gets quickly relativistic.
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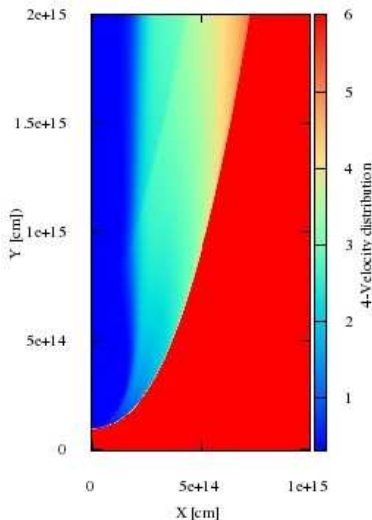
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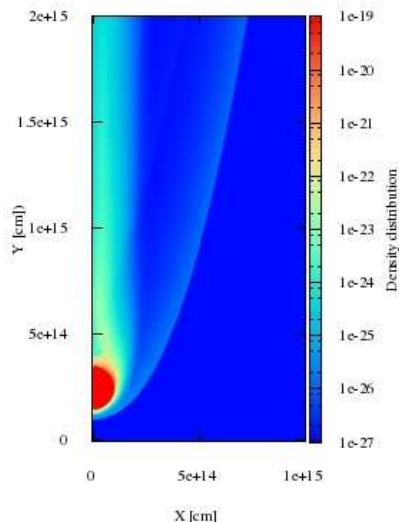
Steady state map of  $\gamma|\beta|$ .



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Steady state density map.

## Radiation: relevant ingredients

- 1 **The power of non-thermal particles comes from the internal energy produced at the interaction.**
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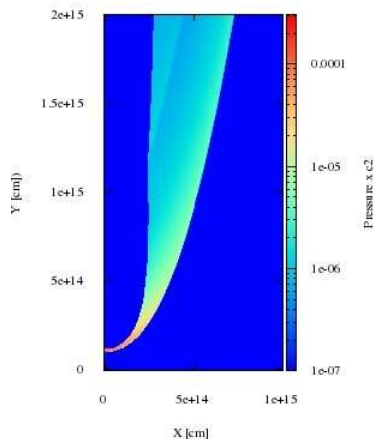
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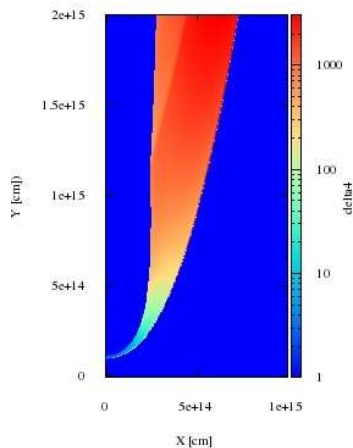
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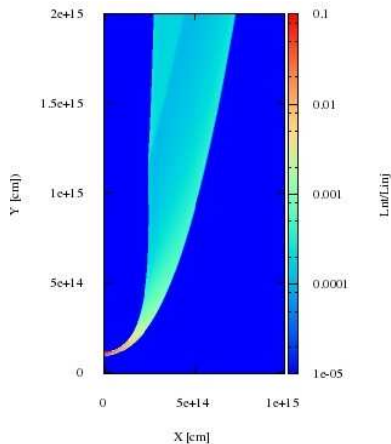


Pressure distribution in the interaction region.

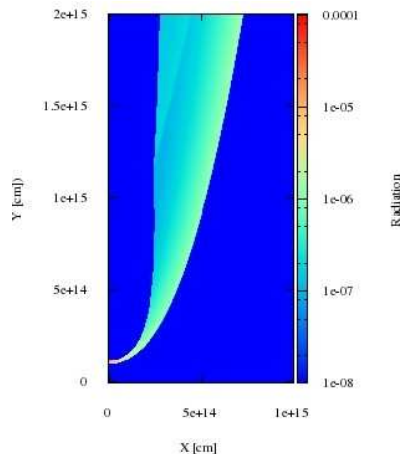


Distribution of  $\delta^4$  (on axis).

# Radiation and dynamics coupling



Radiation efficiency distribution in the interaction region.



Distribution of radiation (on axis).

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# Conclusions

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$$L_{\text{NT}} \lesssim \langle (R_{\text{o,eff}}/R_j)^2 \rangle L_j \sim 10^{37} (R_{\text{o}13}/R_{\text{j,pc}})^2 L_{\text{j}45}$$

- 2 **Ram pressure balance leads to a constant luminosity for stellar winds:**

$$L_{\text{NT}} \lesssim (R_{\text{o,eff}}/R_o)^2 \langle \dot{M}_w v_w \rangle c \sim 10^{37} \dot{M}_{\text{w}25} \text{ erg/s}$$

- 3 **The non-thermal efficiency ( $\eta_{\text{NT}}$ ), the radiation to adiabatic cooling ratio ( $t_{\text{ad/rad}}$ ), and Doppler boosting ( $\delta^4 \Gamma^{-2}$ ), and the distribution of obstacles ( $N_o$ ; clumps, stars or clouds), will determine the global outcome.**

**A persistent (sometimes variable) component of the non-thermal emission could come from jet-obstacle interactions; in the observer frame:**

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- 3 The non-thermal efficiency ( $\eta_{\text{NT}}$ ), the radiation to adiabatic cooling ratio ( $t_{\text{ad/rad}}$ ), and Doppler boosting ( $\delta^4 \Gamma^{-2}$ ), and the distribution of obstacles ( $N_o$ ; clumps, stars or clouds), will determine the global outcome.

A persistent (sometimes variable) component of the non-thermal emission could come from jet-obstacle interactions; in the observer frame:

$$L_{\text{NT,obs}} \sim \langle L_{\text{NT}} \delta^4 \Gamma^{-2} \eta_{\text{NT}} t_{\text{ad/rad}} N_o \rangle \sim 10^{45} L_{\text{NT37}} \delta_{20}^4 \Gamma_{10}^{-2} \eta_{\text{NT-1}} t_{\text{ad/rad,0}} N_{\text{o6}}^{100} \text{ pc erg/s}$$

- 1 Clouds have a fixed radius before disruption. In the lab frame:

$$L_{\text{NT}} \lesssim \langle (R_{\text{o,eff}}/R_j)^2 \rangle L_j \sim 10^{37} (R_{\text{o13}}/R_{j,\text{pc}})^2 L_{j45}$$

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